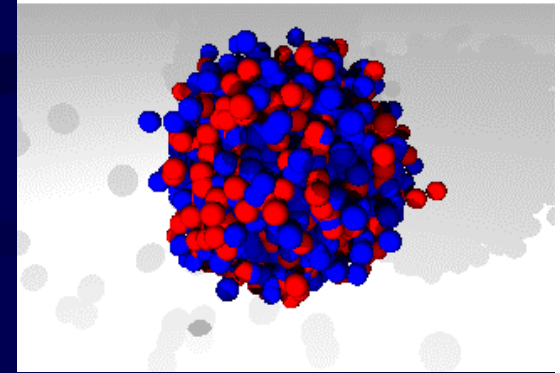
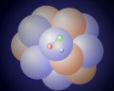


Oliver Buß, Luis Alvarez-Ruso, Pascal Mühlich,
Ulrich Mosel, Radhey Shyam



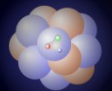
Low-energy pions in nuclear matter

... within a Boltzmann-Uehling-
Uhlenbeck transport model



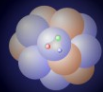
Outline

- Motivation
 - 2π correlations in experiment by TAPS (Metag, Messchendorp, Schadmand et al)
- Part I: Low-energy pions in a BUU transport simulation
 - diploma thesis at JLU Giessen, May 2004
 - http://theorie.physik.uni-giessen.de/html/diploma_theses.html
- Part II: $\pi\pi$ photoproduction off complex nuclei
 - P. Mühlich, L. Alvarez-Ruso, O. Buß, U. Mosel
 - Phys.Lett.B595:216-222, 2004
- Conclusions and Outlook



Motivation

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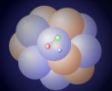
„In-Medium Hadron Physics“

Giessen, November 2004

Motivation:

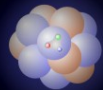
2π correlations with TAPS

- Photoproduction off complex nuclei with 400-460 MeV γ 's at TAPS@MAMI
 - Metag, Messchendorp, Schadmand et al
 - Phys. Rev. Letters 89(2002)222303
- Excitation of the “f0(600)” ($L=0, I=0$) resonance, which is called σ -meson.
 - Possible Final states : $\sigma \rightarrow \pi^0 \pi^0$, $\sigma \rightarrow \pi^+ \pi^-$
 - σ decays already inside the nuclear medium into a $\pi \pi$ final state.
 - Key idea : Observe the invariant mass in the $\pi^0 \pi^0$ channel, to analyze the modification of the σ resonance in the medium.
 - Prediction: The mass of the σ resonance is shifted down
 - e.g Chanfray et al : Phys.Lett.B256 (1991) , Roca et al: Phys.Lett.B541(2002) , ...
- Problem
 - Ordinary final state interactions of pions in the media.



Part I

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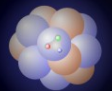
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„In-Medium Hadron Physics“

Giessen, November 2004

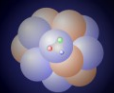
Pions in a Boltzmann-Uehling-Uhlenbeck (BUU) transport simulation



The BUU model

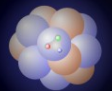
- **Transport model derived in 1930's based on classical Boltzmann equation**
 - Utilized by Bertsch & Kruse in the early 1980's to model heavy ion collisions
 - Phys Rev C29(1984),673
 - Related by limiting procedure to full Kadanov-Baym equations (Quantum mechanical transport equations)
 - Allows full coupled channel calculations
- **Off-Shell transport incorporated in the late 90's**
 - e.g. Effenberger, Leupold, Cassing, Mosel at JLU Giessen (Phys.Rev.C60:051901,1999)
- **Usage of basically the same code in heavy-ion collisions, photon-nucleus and electron-nucleus processes**

$$\frac{df}{dt} = \partial_t f(\vec{p}, \vec{r}) + \nabla_{\vec{p}} H \nabla_{\vec{r}} f(\vec{p}, \vec{r}) - \nabla_{\vec{r}} H \nabla_{\vec{p}} f(\vec{p}, \vec{r}) = I_{coll}$$



Some earlier works...

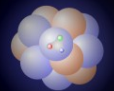
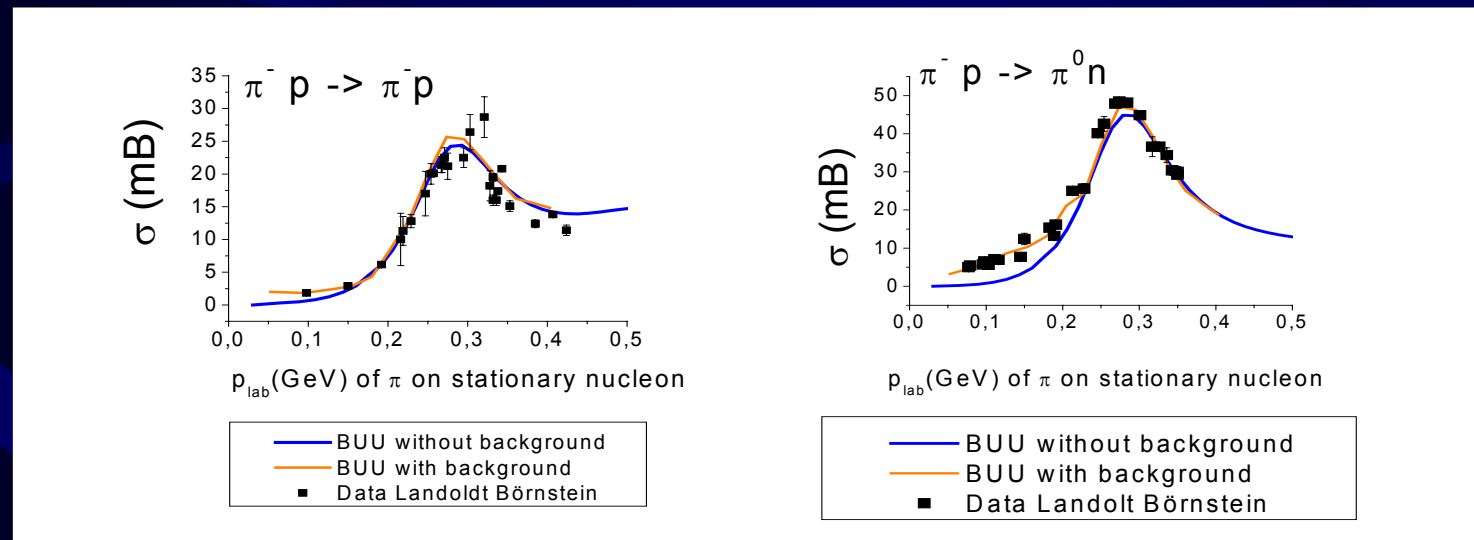
- Salcedo, Oset et al
 - NPA 484 (1988)
 - „Computer simulation of inclusive pion nuclear reactions“
 - $85 \text{ MeV} < E_{\text{kin}} < 300 \text{ MeV}$
 - Only pions as degrees of freedom in simulation
- Engel, Cassing, Mosel, Schäfer and Wolf
 - NPA 572 (1994), 657
 - „Pion nucleus reactions in a microscopic transport model“
 - $85 \text{ MeV} < E_{\text{kin}} < 315 \text{ MeV}$
 - Precursor model of the actual BUU code
- ...



The collisions: Elastic scattering

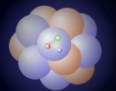
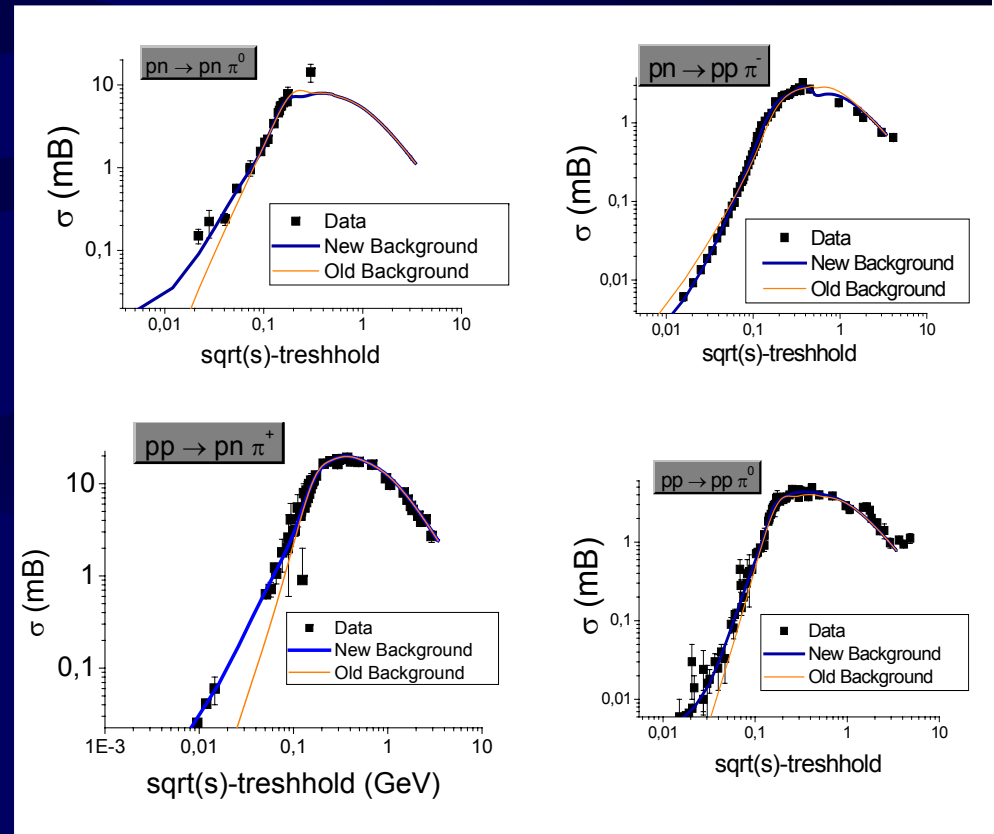
Two channels:

1. $N \pi \rightarrow R \rightarrow N \pi$
2. Non resonant background (not necessary in $I=3/2$ channel)
 - Fitted to match experimental cross sections



The collisions : Pion absorption

- Use detailed balance to describe $N N \pi \Rightarrow N N$ by the inverse process
- Two channels for $N N \pi \Rightarrow N N$:
 1. $N N \Rightarrow N R \Rightarrow N N \pi$
 2. Nonresonant background
 - Fitted to match experimental crosssection for elementary processes



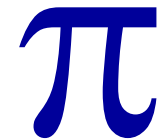
Most important medium modifications in the model

- Pauli blocking of final states if Δ decays
- Delta receives collisional width (new decay channels in medium)
- Potential (Skyrme type with momentum dependence)
- Off-Shell propagation

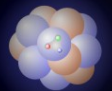
M. Effenberger,
Diploma thesis 1996



- Pauli blocking of final states if π is absorbed
- Potential derived from self energy



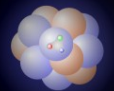
- Potential (Skyrme type with momentum dependence)
- Pauli blocking



The Δ Width

Studied by M. Effenberger, Diploma thesis 1996

- Problem : Missing $\Delta NN \rightarrow NNN$ process.
- Correct collisional width introduced according to Oset et al, NPA 468(1987), 631-652
- Technical: No explicit $\Delta N \rightarrow NN$ process, solved via decay of the Δ 's.



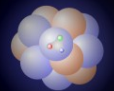
The Δ Potential

Phenomenology: $V_{\Delta} \simeq -30 \text{ MeV}$ (e.g. Ericson-Weise: “Pions and nuclei”, page 248)

Momentum independent nucleon potentials are roughly 50 MeV strong.
Therefore we model the Δ potential by

$$V_{\Delta} = \frac{2}{3} V_{\text{nucleon}}$$

Therefore we also assume same momentum dependence



The Nucleon Potential

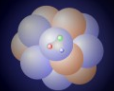
Welke et al, Phys Rev. C38 (1988)

Medium-momentum-dependend potential with $K=290 \text{ MeV}$. Also other options possible!

Fixed in momentum dependence by:

$$\begin{aligned} V_{nucleon}(p = 0, \rho = \rho_0) &= -75 \text{ MeV} \\ V_{nucleon}(p = 0, \rho = \rho_0) &= -75 \text{ MeV} \\ V_{nucleon}(p = \infty, \rho = \rho_0) &= 30.5 \text{ MeV} \end{aligned}$$

N



Effective Hamiltonian for the pion used in BUU

$p_\pi < 140$ MeV:

Low energy model for the self energy Π

(Nieves et al, NPA 554 (1993))

- Valid for low momentum
- Evaluate self energy at vacuum energy
- Good approximation at low energies

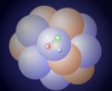
$$H_{\text{eff}} \equiv \sqrt{m^2 + \vec{p}^2} + \Re(\Pi(m^2 + \vec{p}^2)) + V_C$$

$p_\pi \geq 140$ MeV:

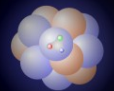
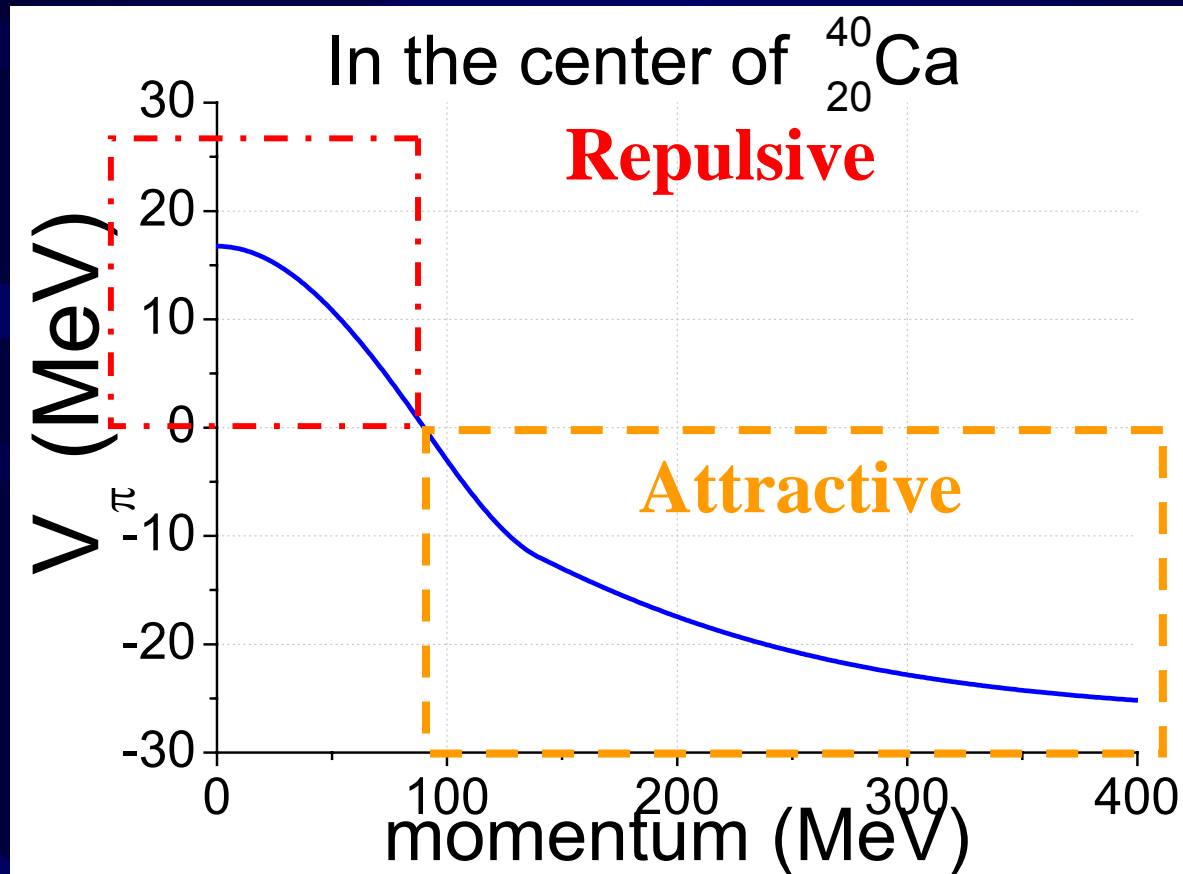
Use simple delta hole model

- Solve dispersion relation analytically $p=p(\omega)$
- Derive spectral function $\rho(\omega, p)$
 - Result: Two branches in $\rho(\omega, p)$
 - W. Eehalt
 - Diploma thesis 1992
 - P. Henning,
 - Phys.Rep. 253 (1995)

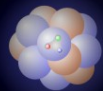
$$H_{\text{eff}} \equiv Z_{\Delta\text{-hole-branch}} E_{\Delta\text{-hole-branch}} + Z_{\pi\text{-branch}} E_{\pi\text{-branch}} + V_C$$



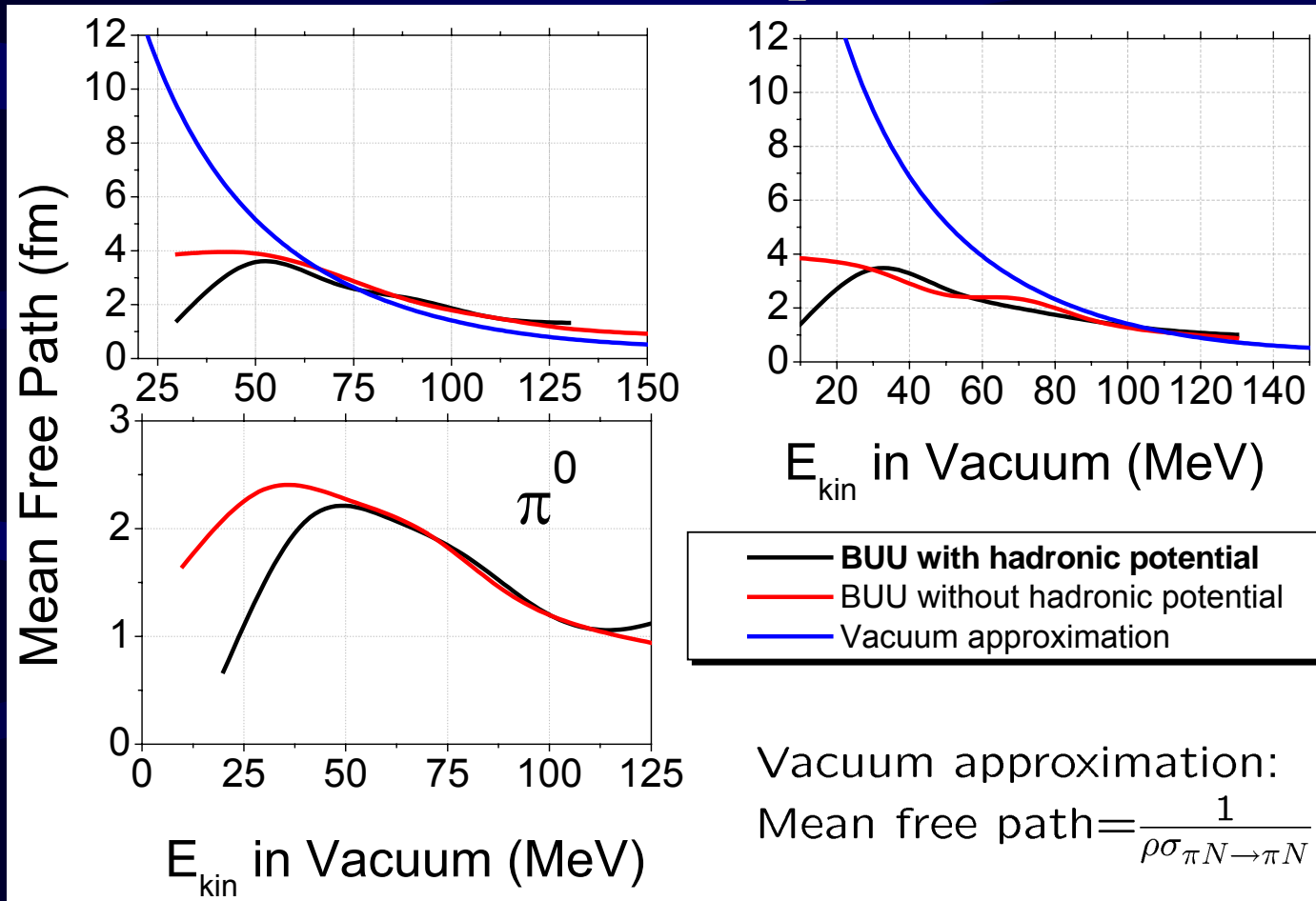
$$V_{\pi} = H_{eff} - H_{vacuum}$$



Results I : Mean free path in BUU

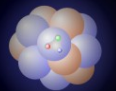


Mean free path

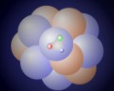
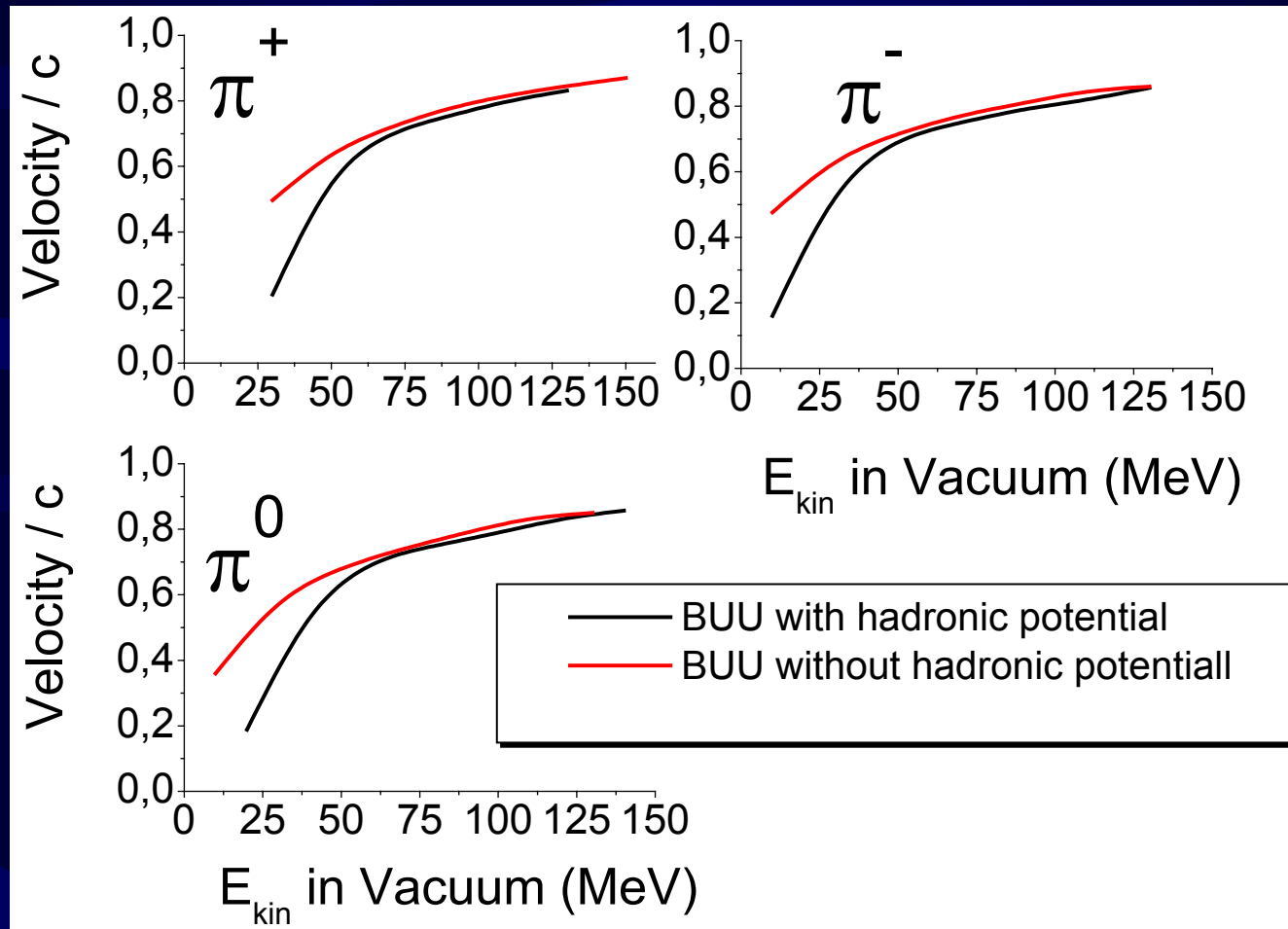


⇒ $NN \leftrightarrow N N \pi$ is important

- Minor impact of potential at large energies
- Steep decrease at low energies → problems with tunneling?!



Velocity modification



Comparison to result with optical model

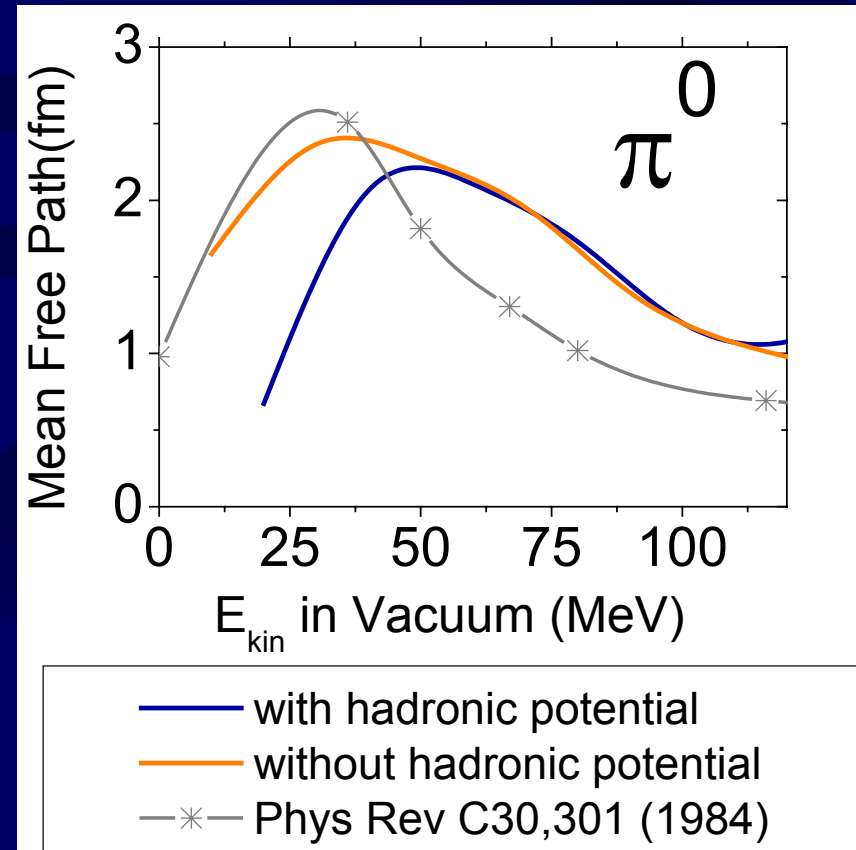
- Optical potential based on analysis of scattering experiments
- Determine mean free path with

$$\Psi_{\pi} = e^{i\mathbf{k}\mathbf{r}} \text{ with } \mathbf{k} \in \mathbb{C}.$$

$$\Rightarrow I(\mathbf{r}) \sim \Psi_{\pi}^* \Psi_{\pi} = e^{-\frac{r}{\lambda}} \text{ with } \lambda = \frac{1}{2\Im(\mathbf{k})}$$

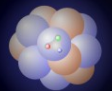
- Determine k by solving K.-G. equation in momentum space

$$(\mathbf{E} + \mathbf{V}_{\text{opt}}(\mathbf{E}, \mathbf{k}))^2 = \mathbf{k}^2 + \mathbf{m}^2$$



Mehrem, Radi, Rasmussen

Phys Rev C30,301 (1984)



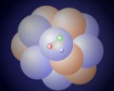
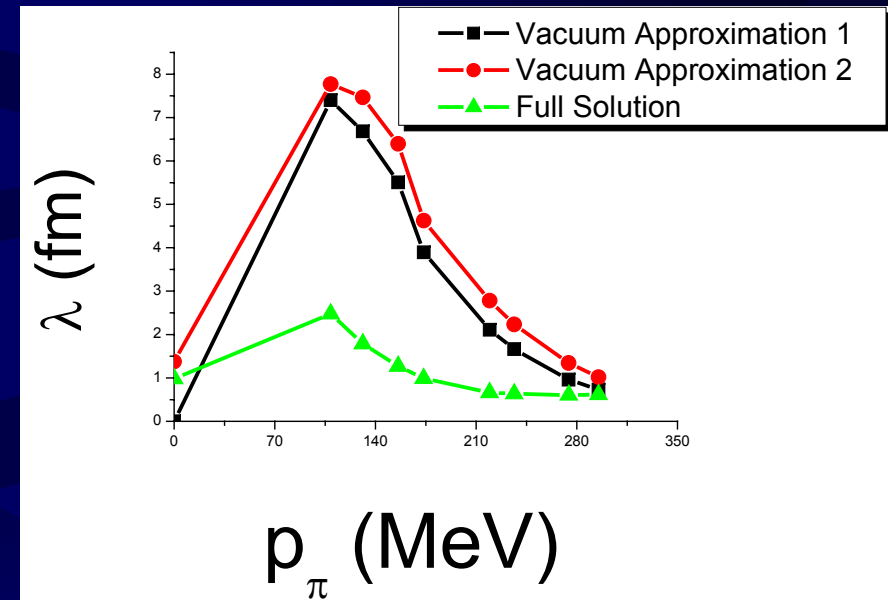
Old knowledge

- Solving the dispersion relation is (in some cases) the key
- Vacuum approximations faulty :

$$\lambda \simeq -\frac{p}{\text{Im} \left(2\sqrt{p^2 + m^2} V_{\text{opt}}(p, \omega = \sqrt{p^2 + m^2}) \right)}$$

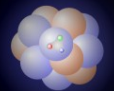
$$\lambda \simeq -\frac{p}{2\sqrt{p^2 + m^2} \text{Im} \left(V_{\text{opt}}(p, \omega = \sqrt{p^2 + m^2}) \right)}$$

- Much too high results when evaluating the mean free path with a given optical potential and the vacuum dispersion relations.

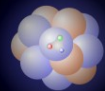


The Message

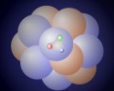
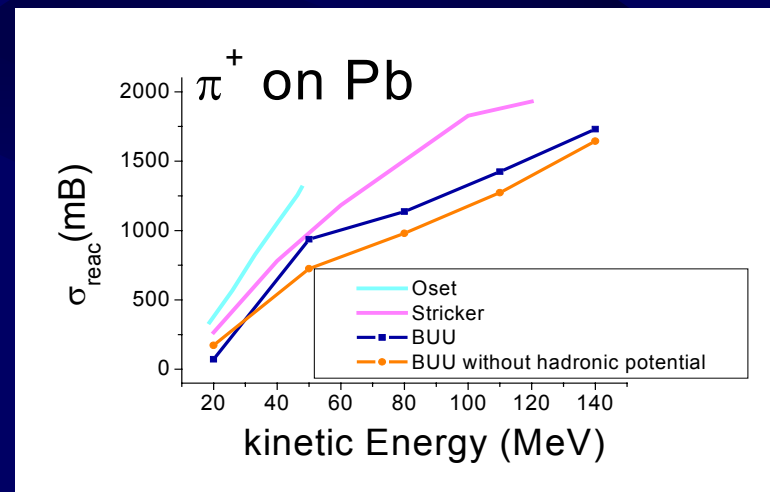
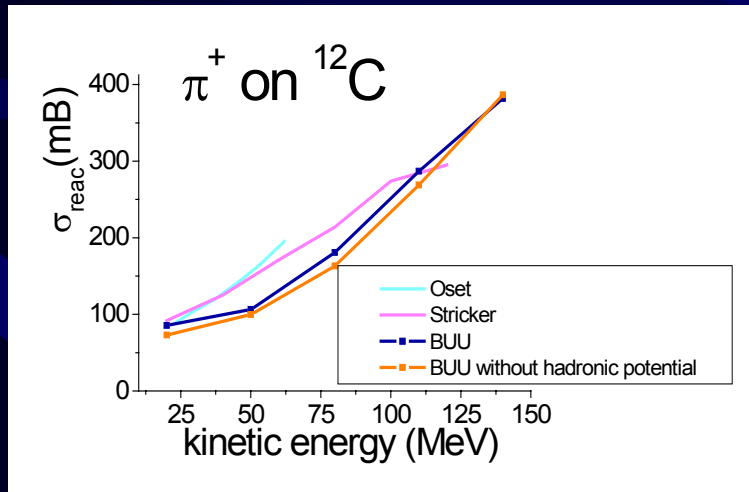
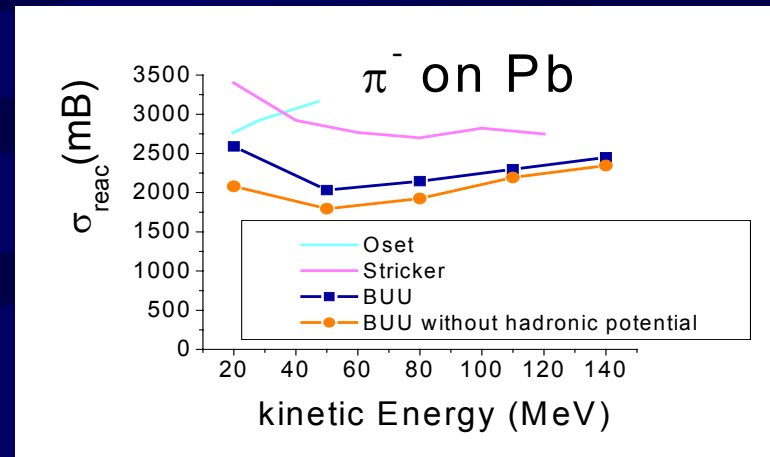
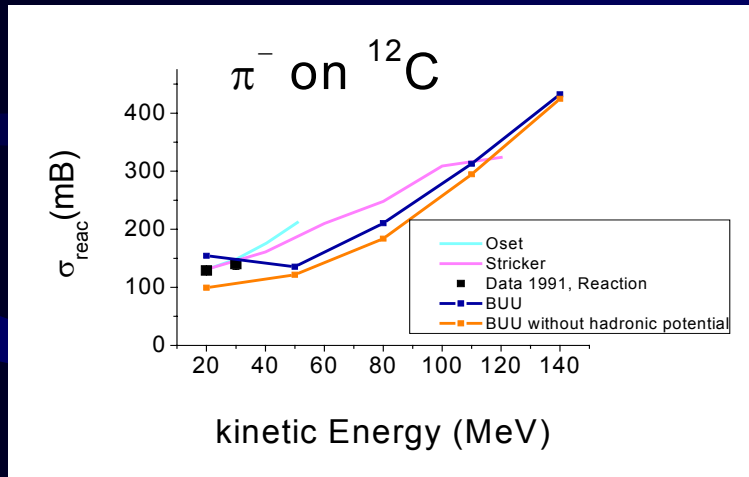
1. Not straight-forward how to define the classical mean free path out of optical models
2. More than one branch possible in microscopic calculations: Averaging ???
3. Microscopic models valid in vacuum approximation ?!



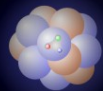
Results II: Pion on nucleus



Quantum vs. BUU

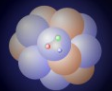
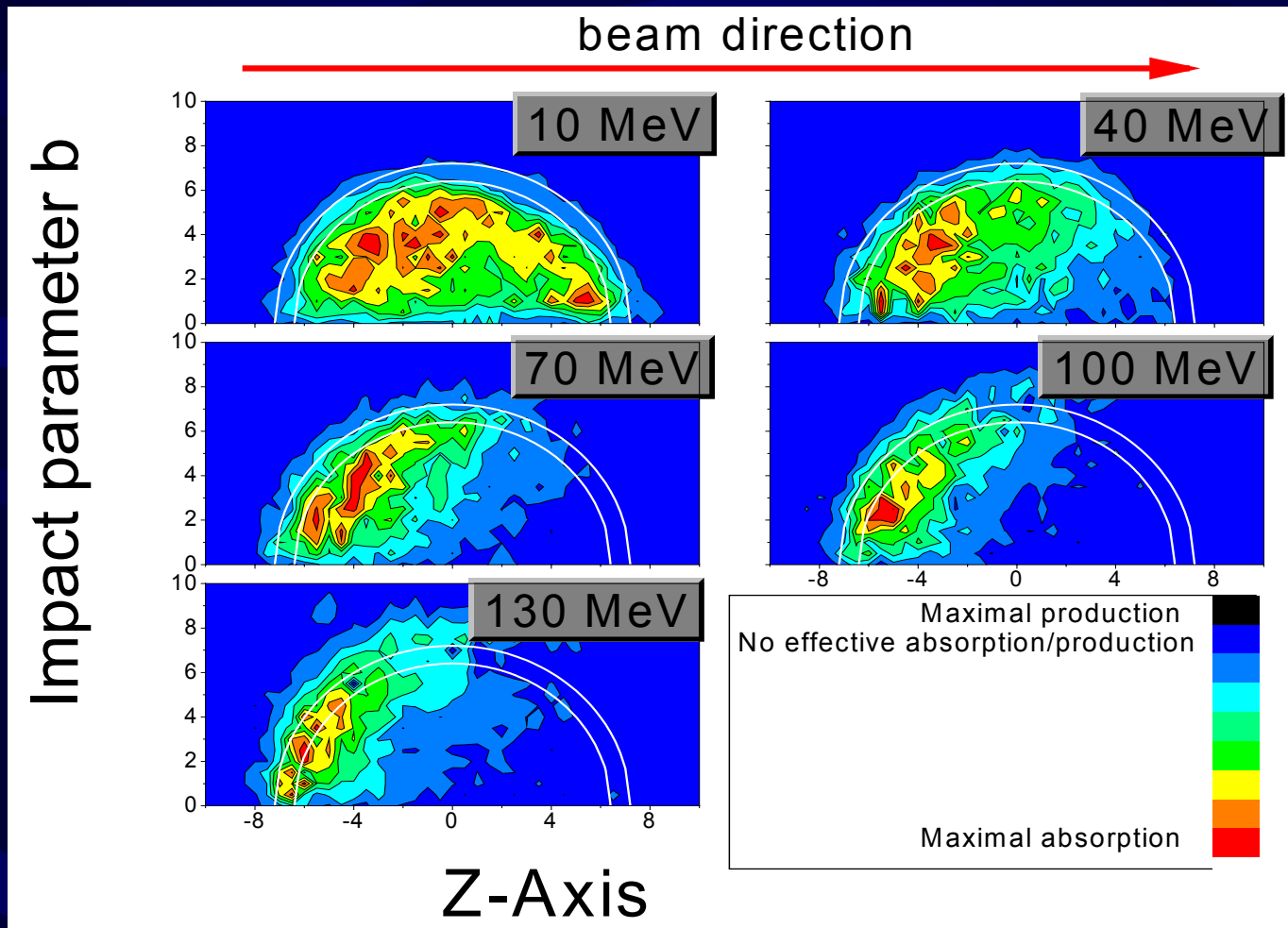


Results III : Pion absorption

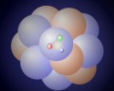
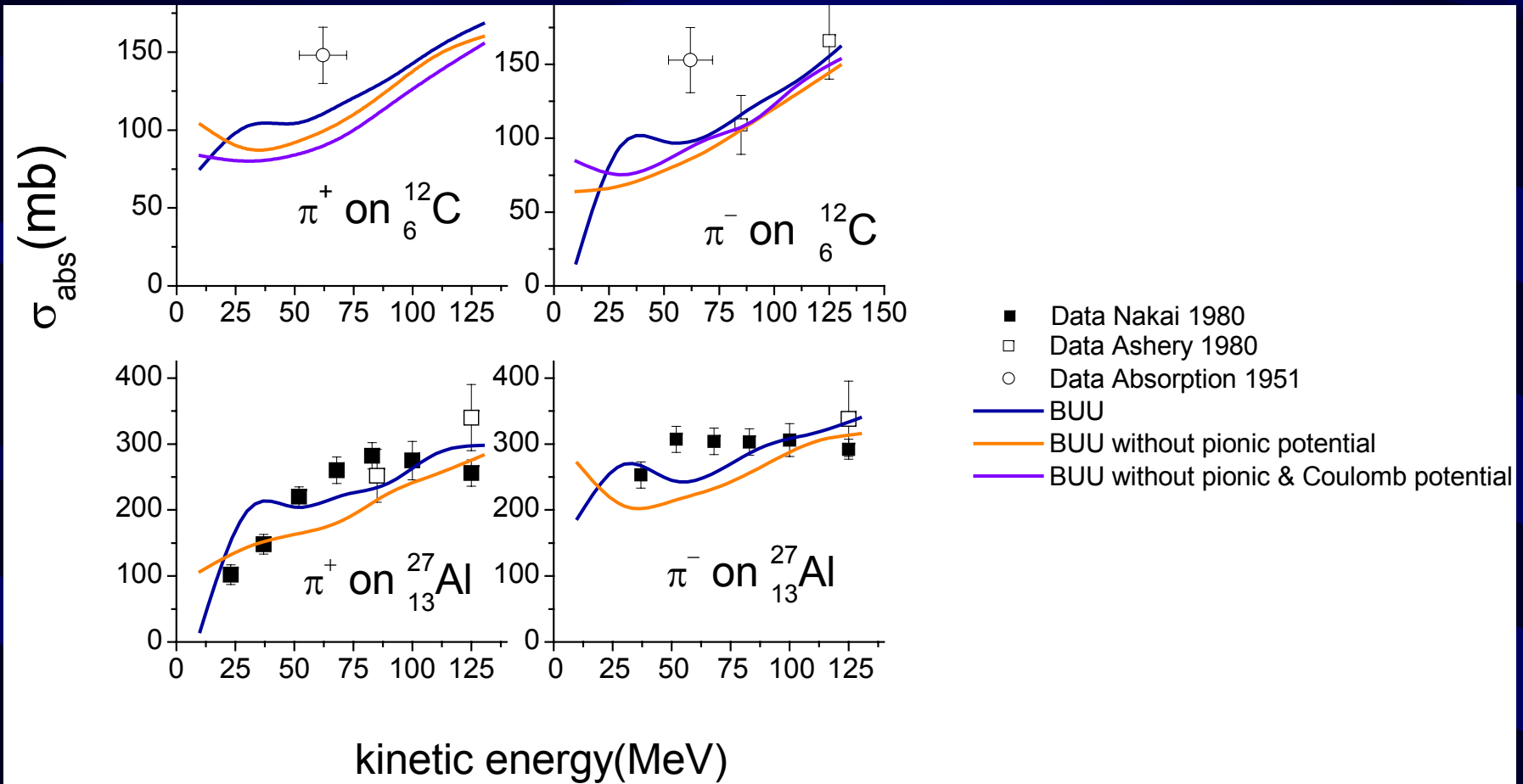


Place of absorption:

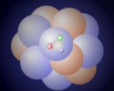
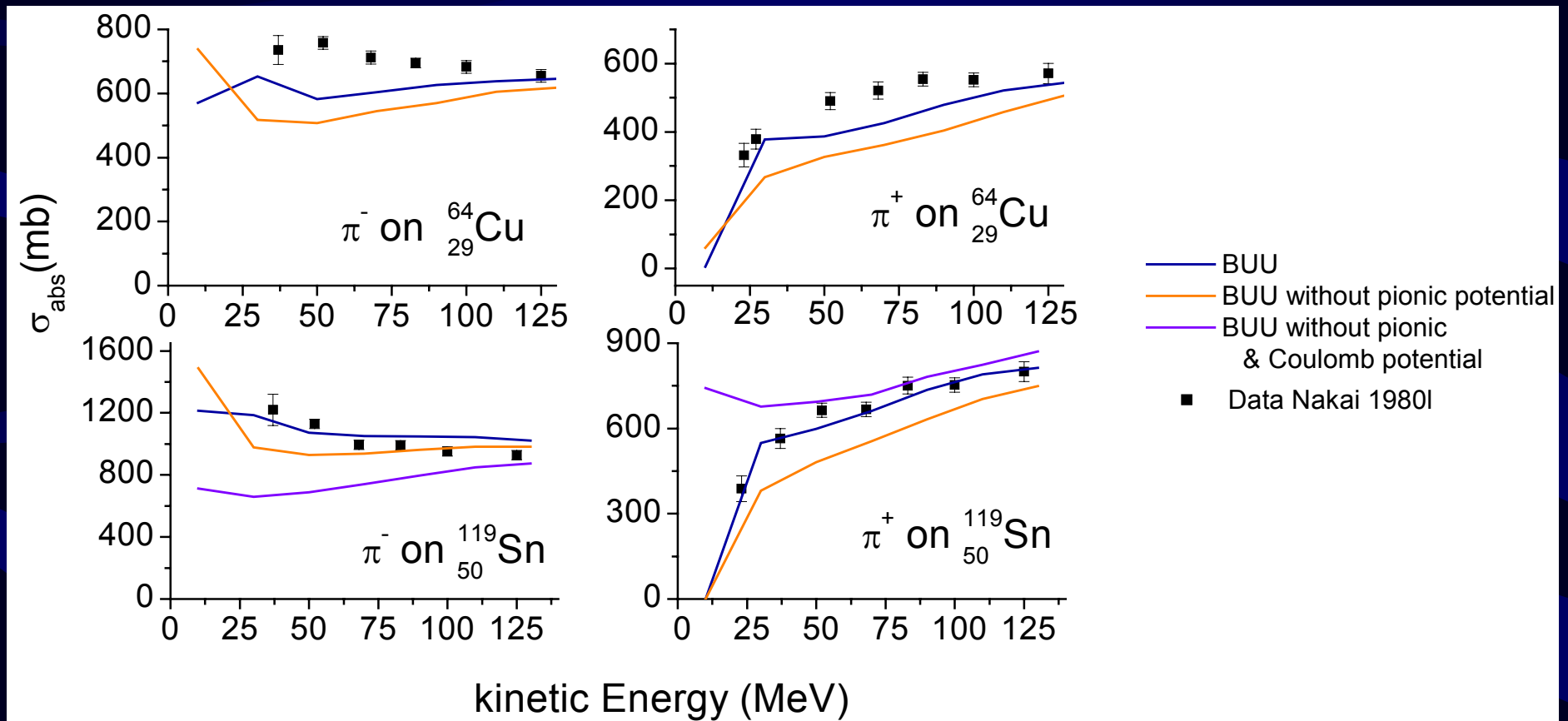
π^- on Pb, without optical potential



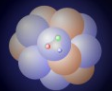
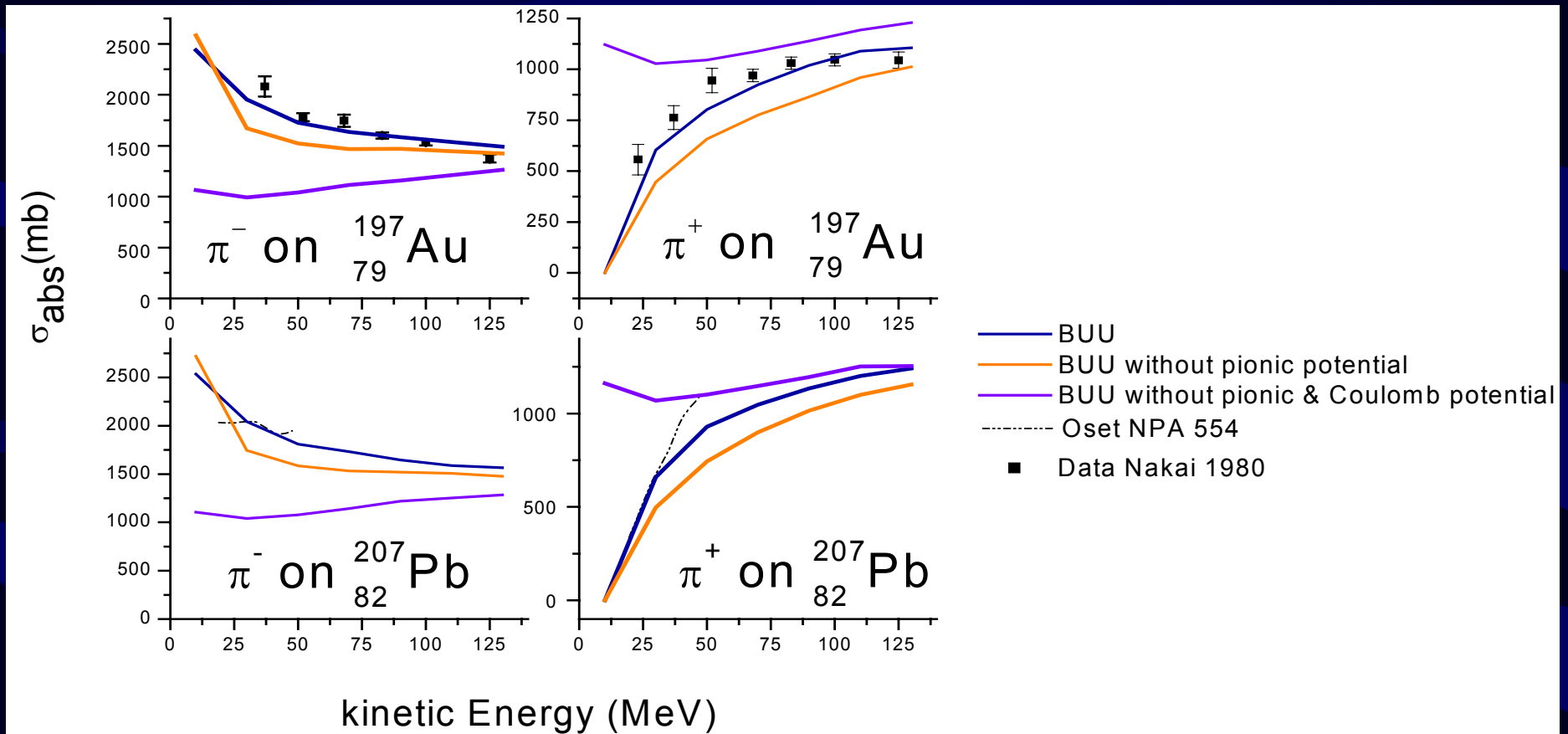
Absorption: BUU – light nuclei



Absorption: BUU –medium size nuclei

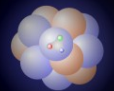


Absorption: BUU – heavy nuclei



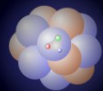
Conclusions Part I

- Sizeable impact of pion potential in absorption experiment
- Absorption is described in a good manner
- Decay width compatible to optical model results
- Important how to define a Mean Free Path
 - Some discrepancies in literature



Part II

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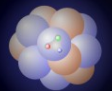
„In-Medium Hadron Physics“

Giessen, November 2004

Photo induced :

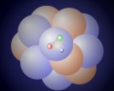
$$\gamma A \rightarrow A' \pi \pi$$

Pascal Mühlich, Luis Alvarez-Ruso, Oliver Buß, U. Mosel



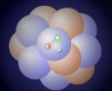
Why is $\gamma A \rightarrow A' \pi \pi$ interesting?

- As intermediate process: $\sigma \rightarrow \pi \pi$
- Prediction of lower σ mass in medium, therefore shift of $d\sigma/dm_{\pi\pi}$ spectra expected.



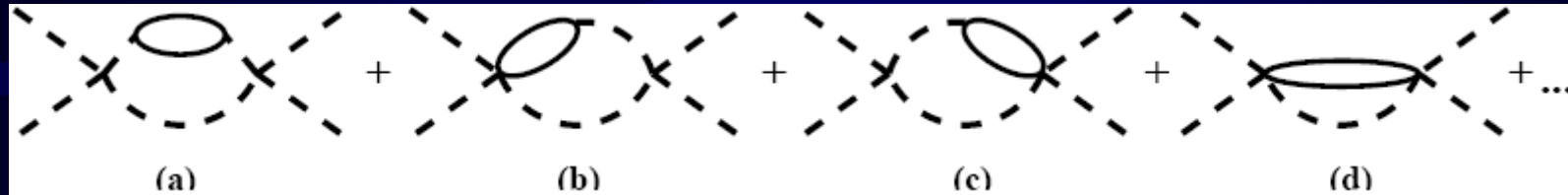
Different models:

- Linear σ model includes σ -meson as chiral partner of the pion
 - Expect lowering of the mass due to chiral symmetry restoration in the medium
- Dynamically generated pole in the $\pi \pi$ scattering amplitude
 - Oller, Oset et al, PRD 59 (1999)
 - In the medium shift to lower masses due to strong p-Wave couplings of pions to ph – and Δh -states.



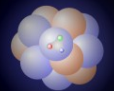
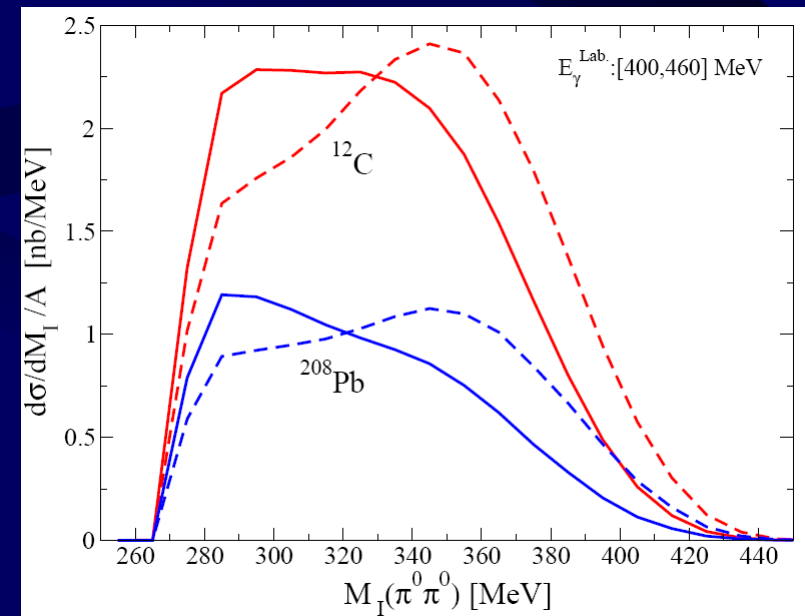
Roca, Oset, Vicente Vacas

Phys.Lett.B541:77-86,2002



Strong particle-hole couplings in the 2π – final rescattering term.

$\Rightarrow \sigma$ gets lighter and more narrow in the medium



Our approach

- Initialization

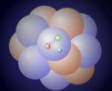
- Elementary 2π photo-production
 - Tejedor, Oset, NPA 600 (96)
 - Nacher et al , NPA 695 (01)
- Vacuum cross sections, no medium modifications in this place

Final State 1

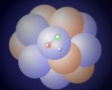
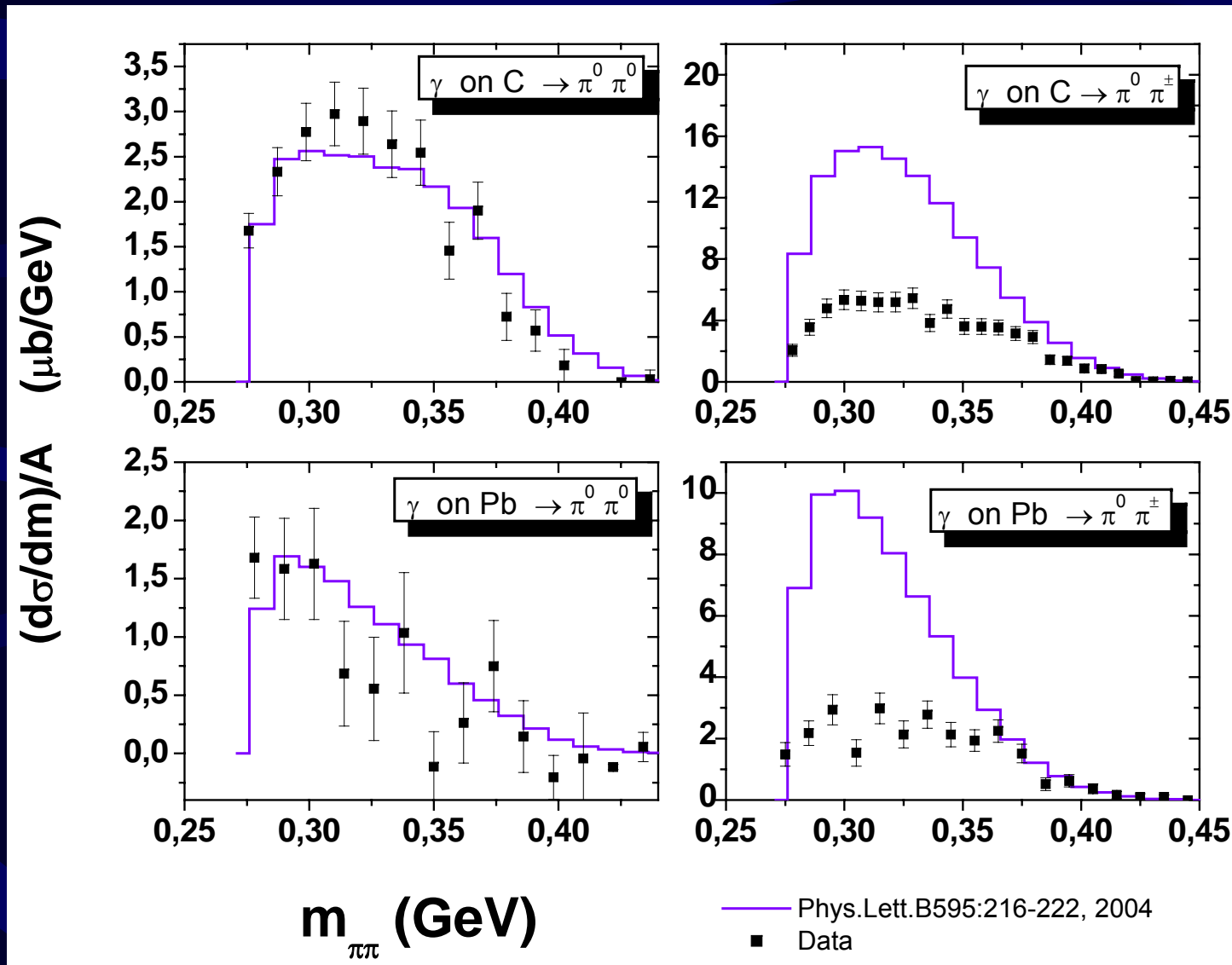
- Standard BUU
- With propagation and medium-modification of the resonances
- Neglect $\pi N \rightarrow \pi \pi N$
- Paper in preparation

Final State 2

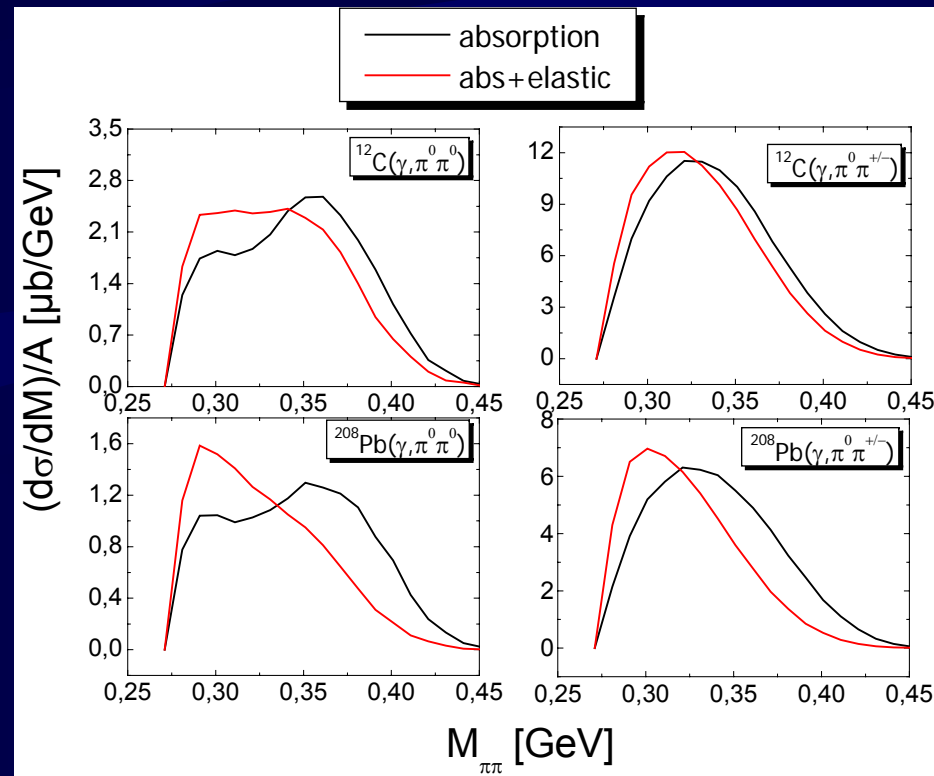
- Absorption by optical model
 - Nieves et al, NPA 554 (1993)
 - Garcia-Recio et al, NPA 526 (1991)
- Elastic scattering according to BUU prescription
- No resonances in the simulation, but the correct angular distribution for πN (SAID)
- Neglect $\pi N \rightarrow \pi \pi N$
- Phys.Lett.B595:216-222, 2004



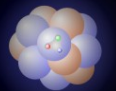
Here: Final State 2.



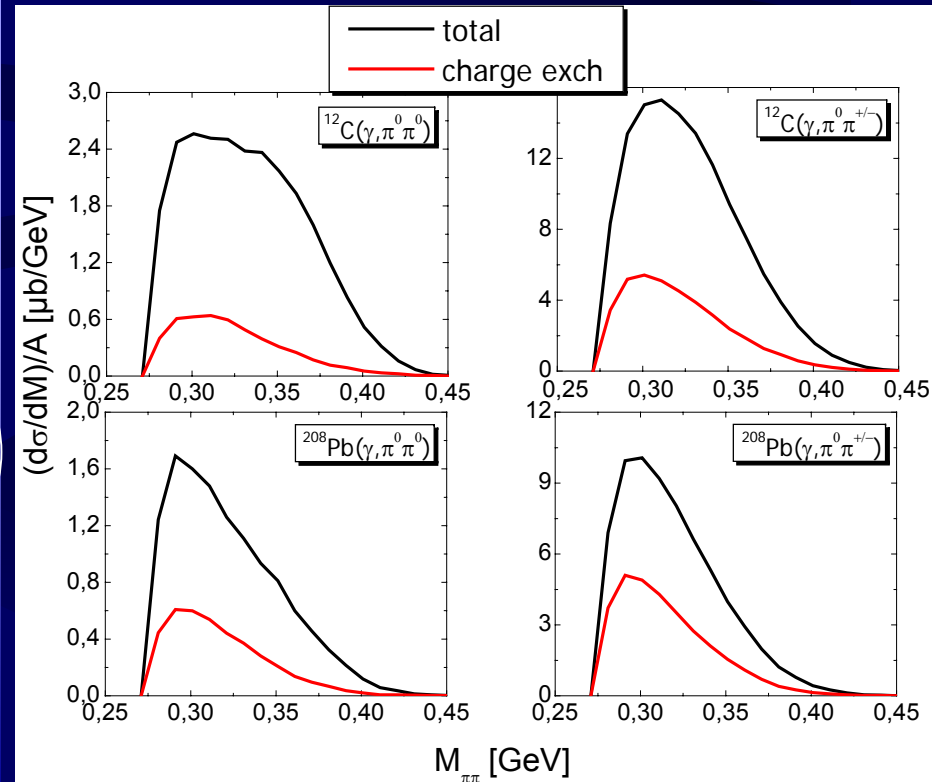
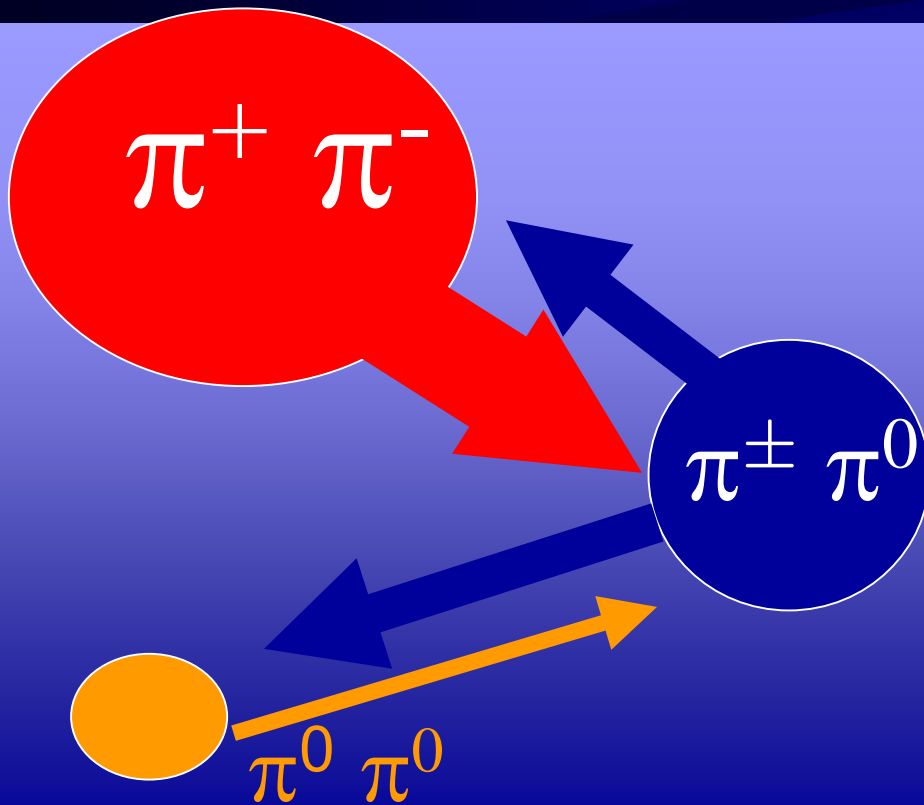
Considering quasielastic scattering in the final state



Message : Not negligible contributions due to elastic scattering. **Here: Final State 2.**

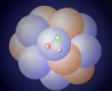


Considering charge exchange in the final state



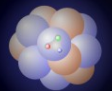
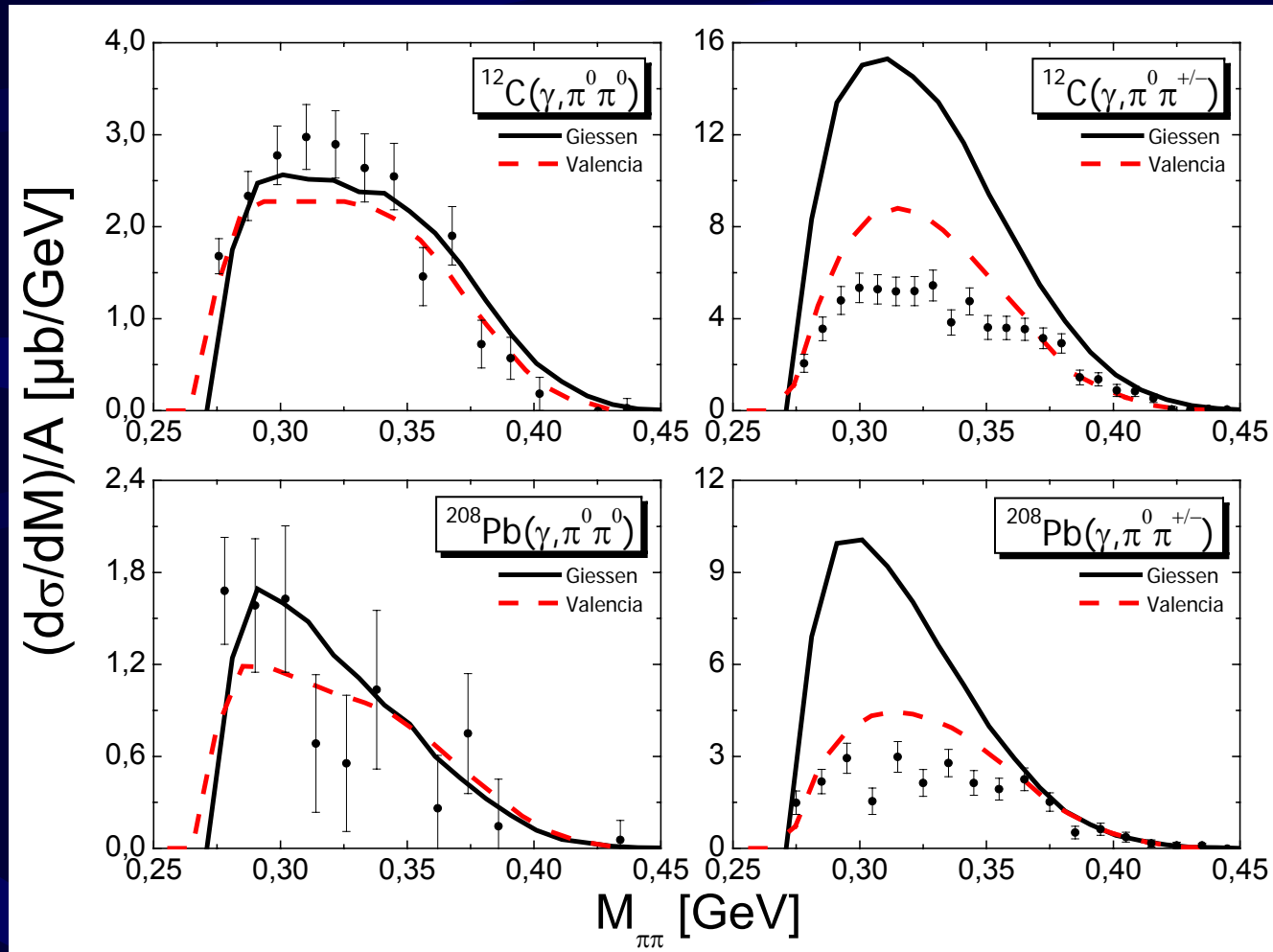
Large contributions due to charge exchange, since different charge channels are not of same size in the initialization.

Here: Final State 2.



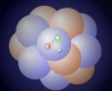
Comparison to Roca et al

Phys.Lett.B541:77-86,2002



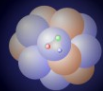
Conclusions Part II

- Importance of ordinary final state effects
- Observed target mass dependence of the $\pi^0 \pi^0$ channel is no unmistakable evidence of the modification of the $\pi\pi$ interaction in the medium
- Discrepancy with the data in the $\pi^0 \pi^\pm$ channel is a standing problem...



Thank you ...

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„In-Medium Hadron Physics“

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